

The equation with Fourier transformation

$$\left(\frac{1}{2} (-i\nabla + \vec{k})^2 + V(\vec{r}) \right) u_{n,\mathbf{k}}(\vec{r}) = E_n(\mathbf{k}) u_{n,\mathbf{k}}(\vec{r})$$

$$V(\vec{r}) = V_{ext}(\vec{r}) + V_H(\vec{r}) + V_{xc}(\vec{r})$$

$$V_{ext}(\vec{r}) = \sum_{\mathbf{R}} \sum_{\tau=1}^{N_{basis}} \frac{-Z_{\tau}}{|\vec{r} - \vec{R} - \vec{\tau}|}$$

$$V_H(\vec{r}) = \int \frac{\rho(\vec{r}')}{|\vec{r} - \vec{r}'|} d^3\vec{r}'$$

$$V_{xc}(\vec{r}) = V_{xc}[\rho(\vec{r})]$$

Fourier transformation of the potential

$$\left(\frac{1}{2} (-i\nabla + \vec{k})^2 + V(\vec{r}) \right) u_{n,\mathbf{k}}(\vec{r}) = E_n(\mathbf{k}) u_{n,\mathbf{k}}(\vec{r})$$

$$V(\vec{r}) = \sum_l V(\vec{G}_l) e^{i\vec{G}_l \cdot \vec{r}} \quad , \quad V(\vec{G}) = \frac{1}{\Omega_{cell}} \int_{cell} V(\vec{r}) e^{-i\vec{G} \cdot \vec{r}} d^3\vec{r}$$

Fourier transformation of the potential

$$\left(\frac{1}{2}(-i\nabla + \vec{k})^2 + V(\vec{r}) \right) u_{n,\mathbf{k}}(\vec{r}) = E_n(\mathbf{k}) u_{n,\mathbf{k}}(\vec{r}) \quad , \quad u_{n,\mathbf{k}}(\vec{r}) = \sum_{p=1}^{\infty} C_p e^{i\vec{G}_p \cdot \vec{r}}$$

$$\begin{bmatrix} H_{11} & H_{12} & H_{13} & \dots \\ H_{21} & H_{22} & H_{23} & \dots \\ H_{31} & H_{32} & H_{33} & \dots \\ \dots & \dots & \dots & \dots \end{bmatrix} \begin{pmatrix} C_1 \\ C_2 \\ C_3 \\ \dots \end{pmatrix} = E_n(\mathbf{k}) \begin{pmatrix} C_1 \\ C_2 \\ C_3 \\ \dots \end{pmatrix}$$

$\infty \times \infty$ matrix

The Hamiltonian matrix element

$$H_{m,n} = \left\langle e^{i\vec{G}_m \cdot \vec{r}} \left| \frac{1}{2} (-i\nabla + \vec{k})^2 \right| e^{i\vec{G}_n \cdot \vec{r}} \right\rangle_{\infty} = \frac{1}{2} (\vec{G}_n + \vec{k})^2 \delta_{m,n} + V(\vec{G}_m - \vec{G}_n)$$

The Hamiltonian matrix element

$$\left(\frac{1}{2} (-i\nabla + \vec{k})^2 + \sum_{\mathbf{R}} \sum_{\tau=1}^{N_{basis}} \frac{-Z_{\tau}}{|\vec{r} - \vec{R} - \vec{\tau}|} + \int \frac{\rho(\vec{r}')}{|\vec{r} - \vec{r}'|} d^3\vec{r}' + V_{xc}[\rho(\vec{r})] \right) u_{n,\mathbf{k}}(\vec{r}) = E_n(\mathbf{k}) u_{n,\mathbf{k}}(\vec{r})$$

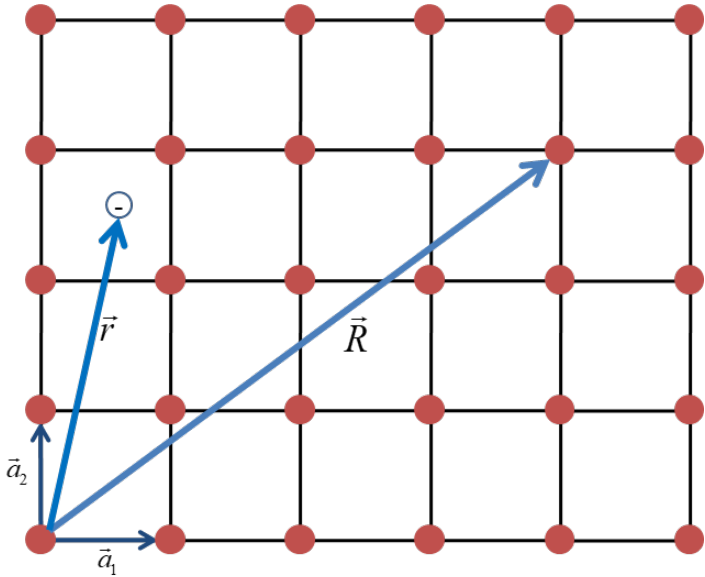
$$\begin{aligned} H_{m,n} &\rightarrow \left\langle e^{i\vec{G}_m \cdot \vec{r}} \left| \sum_{\mathbf{R}} \sum_{\tau=1}^{N_{basis}} \frac{-Z_{\tau}}{|\vec{r} - \vec{R} - \vec{\tau}|} \right| e^{i\vec{G}_n \cdot \vec{r}} \right\rangle_{\infty} = \sum_{\mathbf{R}} \sum_{\tau=1}^{N_{basis}} \int_{\infty} e^{i(\vec{G}_n - \vec{G}_m) \cdot \vec{r}} \frac{-Z_{\tau}}{|\vec{r} - \vec{R} - \vec{\tau}|} = \sum_{\mathbf{R}} \sum_{\tau=1}^{N_{basis}} e^{i(\vec{G}_n - \vec{G}_m) \cdot (\vec{R} + \vec{\tau})} \int_{\infty} e^{i(\vec{G}_n - \vec{G}_m) \cdot \vec{r}} \frac{-Z_{\tau}}{|\vec{r}|} \\ &= \sum_{\mathbf{R}} \sum_{\tau=1}^{N_{basis}} e^{i(\vec{G}_n - \vec{G}_m) \cdot \vec{\tau}} \int_{\infty} e^{i(\vec{G}_n - \vec{G}_m) \cdot \vec{r}} \frac{-Z_{\tau}}{|\vec{r}|} = N_{cell} \sum_{\tau=1}^{N_{basis}} e^{i(\vec{G}_n - \vec{G}_m) \cdot \vec{\tau}} \int_{\infty} e^{i(\vec{G}_n - \vec{G}_m) \cdot \vec{r}} \frac{-Z_{\tau}}{|\vec{r}|} \\ &= -N_{cell} \sum_{\tau=1}^{N_{basis}} e^{i(\vec{G}_n - \vec{G}_m) \cdot \vec{\tau}} \frac{-Z_{\tau} 4\pi}{|\vec{G}_n - \vec{G}_m|^2} \end{aligned}$$

The Hamiltonian matrix element

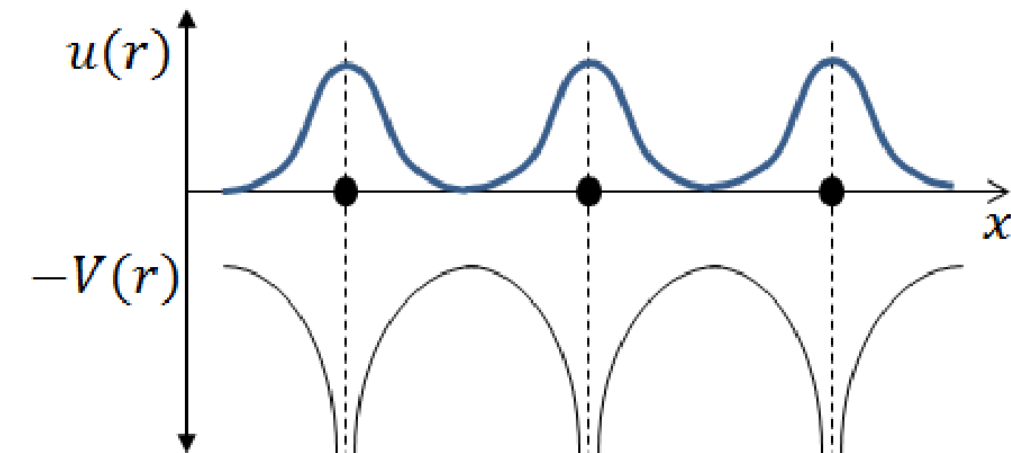
$$\left(\frac{1}{2}(-i\nabla + \vec{k})^2 + \sum_{\mathbf{R}} \sum_{\tau=1}^{N_{basis}} \frac{-Z_{\tau}}{|\vec{r} - \vec{R} - \vec{\tau}|} + \int \frac{\rho(\vec{r}')}{|\vec{r} - \vec{r}'|} d^3\vec{r}' + V_{xc}[\rho(\vec{r})] \right) u_{n,\mathbf{k}}(\vec{r}) = E_n(\mathbf{k}) u_{n,\mathbf{k}}(\vec{r})$$

$$\begin{aligned} H_{m,n} &\rightarrow \left\langle e^{i\vec{G}_m \cdot \vec{r}} \left| \int \frac{\rho(\vec{r}')}{|\vec{r} - \vec{r}'|} d^3\vec{r}' \right| e^{i\vec{G}_n \cdot \vec{r}} \right\rangle_{\infty} = \int_{\infty} e^{i(\vec{G}_n - \vec{G}_m) \cdot \vec{r}} \int \frac{\rho(\vec{r}')}{|\vec{r} - \vec{r}'|} d^3\vec{r}' = \sum_{\vec{G}} \iint e^{i(\vec{G}_n - \vec{G}_m) \cdot \vec{r}} \frac{\rho(\vec{G}) e^{i\vec{G} \cdot \vec{r}'}}{|\vec{r} - \vec{r}'|} d^3\vec{r}' d^3\vec{r} \\ &= \sum_{\vec{G}} \rho(\vec{G}) \int d^3\vec{r} e^{i(\vec{G}_n - \vec{G}_m) \cdot \vec{r}} e^{i\vec{G} \cdot \vec{r}} \int \frac{e^{i\vec{G} \cdot (\vec{r}' - \vec{r})}}{|\vec{r}' - \vec{r}|} d^3\vec{r}' \\ &= \sum_{\vec{G}} \rho(\vec{G}) \frac{4\pi}{G^2} \int d^3\vec{r} e^{i(\vec{G}_n - \vec{G}_m) \cdot \vec{r}} e^{i\vec{G} \cdot \vec{r}} = \sum_{\vec{G}} \rho(\vec{G}) \frac{4\pi}{G^2} N_{cell} \delta_{\vec{G}_n - \vec{G}_m, \vec{G}} = N_{cell} \frac{4\pi \rho(\vec{G})}{|\vec{G}_n - \vec{G}_m|^2} \end{aligned}$$

Hamiltonian with a translational symmetry



Translation Operator



➤ Now let us think of Hamiltonian with a periodic potential

$$V(\vec{r} + \vec{R}) = V(\vec{r})$$

$$\hat{H} = -\frac{1}{2} \nabla^2 + V(\vec{r})$$

$$\hat{H}(\vec{r} + \vec{R}) = \hat{H}(\vec{r})$$