# Hamiltonian with a translational symmetry



**Translation Operator** 



Now let us think of Hamiltonian with a periodic potential

$$\begin{split} \hat{V}(\vec{r} + \vec{R}) &= V(\vec{r}) \\ \hat{H} &= -\frac{1}{2}\nabla^2 + V(\vec{r}) \\ \hat{H}(\vec{r} + \vec{R}) &= \hat{H}(\vec{r}) \end{split}$$

## **Statements for symmetry**

- I. The Hamiltonian is invariant over the translation
- II. The Hamiltonian commutes with the translation operator.

$$\left[\hat{H},\hat{T}(\vec{R})\right] = 0 \quad ,\hat{H}\hat{T} - \hat{T}\hat{H} = 0$$

$$\hat{T}^{-1}\hat{H}\hat{T} = \hat{H}$$

$$\hat{T}(\vec{R})\hat{H}(\vec{r})\psi(\vec{r}) = \hat{H}(\vec{r}+\vec{R})\psi(\vec{r}+\vec{R}) = \hat{H}(\vec{r})\psi(\vec{r}+\vec{R}) = \hat{H}\hat{T}(\vec{R})\psi(\vec{r})$$

# Good quantum number, stationary state

- I. Once the operator commutes with the Hamiltonian, the eigenvalue of the operator can stay there stationary.
- II. Say, at t = 0, we have the eigenstates for translation,  $\hat{T}[\vec{R}]$ , it keeps the same eigenvalue at t > 0.

$$\hat{T}(\vec{R}) \left| \psi_{\vec{k}}(t=0) \right\rangle = e^{i\vec{k}\cdot\vec{R}} \left| \psi_{\vec{k}}(t=0) \right\rangle$$
$$\left| \psi_{\vec{k}}(t) \right\rangle = \exp\left[-i\frac{\hat{H}}{\hbar}t\right] \left| \psi_{\vec{k}}(t=0) \right\rangle$$

$$\hat{T}(\vec{R}) \left| \psi_{\vec{k}}(t) \right\rangle = \hat{T}(\vec{R}) \exp\left[-i\frac{\hat{H}}{\hbar}t\right] \left| \psi_{\vec{k}}(t=0) \right\rangle = \exp\left[-i\frac{\hat{H}}{\hbar}t\right] \hat{T}(\vec{R}) \left| \psi_{\vec{k}}(t=0) \right\rangle = \exp\left[-i\frac{\hat{H}}{\hbar}t\right] e^{i\vec{k}\cdot\vec{R}} \left| \psi_{\vec{k}}(t=0) \right\rangle = e^{i\vec{k}\cdot\vec{R}} \exp\left[-i\frac{\hat{H}}{\hbar}t\right] \left| \psi_{\vec{k}}(t=0) \right\rangle$$
We have  $\hat{T}(\vec{R}) \left| \psi_{\vec{k}}(t) \right\rangle = e^{i\vec{k}\cdot\vec{R}} \exp\left[-i\frac{\hat{H}}{\hbar}t\right] \left| \psi_{\vec{k}}(t=0) \right\rangle$ 
Therefore, it is ...

## Good quantum number, stationary state

- I. To analyze the Hamiltonian spectrum, when the Hamiltonian has a symmetry, it is much convenient to search for the simultaneous eigenstates.
- II. For simultaneous eigenstates for  $\hat{T}[\vec{R}]$  with  $\hat{H}$ ,

$$\hat{H} = -\frac{1}{2}\nabla^2 + V(\vec{r})$$
$$\left(-\frac{1}{2}\nabla^2 + V(\vec{r})\right)e^{i\vec{k}\cdot\vec{r}}u(\vec{r}) = Ee^{i\vec{k}\cdot\vec{r}}u(\vec{r})$$
$$\left(\frac{1}{2}(-i\nabla + \vec{k})^2 + V(\vec{r})\right)u(\vec{r}) = Eu(\vec{r})$$

### Differential equation eigenvalue problem.

I. We want to determine the eigenvalue and the eigenfunction.

$$\left(\frac{1}{2}(-i\nabla+\vec{k})^2+V(\vec{r})\right)u_{\vec{k}}(\vec{r})=E(\vec{k})u_{\vec{k}}(\vec{r})$$

II. The label for the discrete eigenspectral

$$\left(\frac{1}{2}(-i\nabla + \vec{k})^2 + V(\vec{r})\right)u_{n,\vec{k}}(\vec{r}) = E_n(\vec{k})u_{n,\vec{k}}(\vec{r})$$

III. Look up the math text : the operator and BC

#### Differential equation eigenvalue problem.

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$$\left(\frac{1}{2}(-i\nabla + \vec{k})^2 + V(\vec{r})\right) u_{n,\vec{k}}(\vec{r}) = E_n(\vec{k})u_{n,\vec{k}}(\vec{r})$$
$$u_{n,\vec{k}}(\vec{r} + \vec{R}) = u_{n,\vec{k}}(\vec{r})$$

# Question

Suppose we have a cell-periodic charge density, say,  $\rho(\vec{r}) = \rho(\vec{r} + \vec{R})$ . Does the Coulomb potential have the same periodicity ?

$$V_H(\vec{r}) = \int \frac{\rho(\vec{r}')}{|\vec{r} - \vec{r}'|} d^3 \vec{r}'$$

$$V_{H}(\vec{r}) = V_{H}(\vec{r} + \vec{R})$$
 ???